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Powers ontology and the quantum revolution

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Abstract

An Aristotelian philosophy of nature rejects the modern prejudice in favor of the microscopic, a rejection that is crucial if we are to penetrate the mysteries of the quantum world. I defend an Aristotelian model by drawing on both quantum chemistry and recent work on the measurement problem. By building on the work of Hans Primas, using the distinction between quantum and classical properties that emerges in quantum chemistry at the thermodynamic or continuum limit, I develop a new version of the Copenhagen interpretation, a version that is realist, holistic, and hylomorphic in character, allowing for the attribution of fundamental causal powers to human observers and their instruments. I conclude with a critique of non-hylomorphic theories of primitive ontology, including Bohmian mechanics, Everettianism, and GRW mass-density.

Keywords Quantum mechanics · Powers ontology · Causal powers · Aristotelianism · Neo-Aristotelianism · Hylomorphism · Measurement problem · Neo-Humeanism · Quantum chemistry · Thermodynamics · Many-worlds interpretation · Bohmian mechanics · GRW

1 Introduction

Widespread dissatisfaction with Humean and Neo-Humean projects has led to a revival of interest in Aristotle-inspired theories of causal powers. This revival has great potential to illuminate issues in the philosophy of science and of nature. In particular, an Aristotelian perspective on the import of the quantum revolution would open up new avenues of thought. In this paper, I will sketch one such perspective.

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25 In the first section, I describe the basic elements of a powers ontology, in con-
26 trast to its principal competitors, and I propose that there two distinct philosophies
27 of nature correspond to two of these ontologies (Aristotelian and Humean). Then,
28 in Section 2, I argue that the quantum revolution has taken science in the direction
29 of an Aristotelian metaphysics and philosophy of nature, a fact that has been noted
30 by some (including Planck and Heisenberg) but which has not yet been widely rec-
31 ognized in contemporary philosophy of science. This new direction includes three
32 components: potentiality, processes, and (most importantly) the need for a fundamen-
33 tally real domain (beyond the microphysical) that includes experimenters and their
34 instruments.

35 I explain in Sections 3, 4, and 5 why the Aristotelian philosophy of science offers
36 an alternative to the reduction of special sciences to microphysics. An Aristotelian
37 philosophy of nature rejects the modern prejudice in favor of the microscopic, a
38 rejection that is crucial if we are to penetrate the mysteries of the quantum world.

39 The remainder of the paper is a defense of the Aristotelian model that draws on two
40 areas of contemporary science: quantum chemistry and thermodynamics (Section 6)
41 and the measurement problem (Section 7). I argue that the distinction between com-
42 muting (quantal) and non-commuting (classical) properties in quantum theory (a
43 distinction that appears only when models are taken to the thermodynamic or contin-
44 uum limit) provides the basis for a new version of the Copenhagen interpretation, an
45 interpretation that is realist, holistic, and hylomorphic in character. This new version
46 allows for the attribution of fundamental causal powers (both active and passive) to
47 meso- and macro-scopic entities, including human observers and their instruments.

48 My project has encompasses three phases, three goals—of increasingly ambitious
49 character.

- 50 1. Phase 1: sketch a hylomorphic, powerist interpretation of modern quantum
51 theory, arguing that it represents a genuine and stable location in logical space.
- 52 2. Phase 2: argue that there is no empirical evidence against the hylomorphic
53 interpretation—that it is at least as well supported by data and scientific practice
54 as is the microphysicalist, modern alternative.
- 55 3. Phase 3: argue that the empirical evidence supports the hylomorphic interpre-
56 tation over the other alternatives, including old Copenhagen, Bohm, objective-
57 collapse, and Everett interpretations.

58 I will argue for Phase 1 in Sections 4 and 5, and for phase 2 in the Section 6, with
59 special consideration of quantum theories of chemistry and thermodynamics. I'll take
60 up the case for Phase 3 in the concluding Section 7.

61 **2 Four metaphysical options and two philosophies of nature**

62 There is a natural class of phenomena that at least appears to involve a sort of physical
63 or natural modality. This class includes three sub-classes: subjunctive and counter-
64 factual conditionals, dispositions and causal powers, and causal laws of nature (see
65 Koons and Pickavance2017). It would be quite surprising if all three sub-classes

included metaphysically fundamental facts, since it seems that some can be defined by or grounded in the others. Consequently, there are four ontological options:

1. Powerism. Causal powers and dispositions are fundamental.
2. Hypotheticalism. Facts expressed by means of subjunctive conditionals are fundamental.
3. Nomism. Causal laws of nature are fundamental.
4. Neo-Humeanism. None of these are fundamental, but all are grounded in the *Humean mosaic* of categorical qualities distributed across spacetime.

Hypotheticalism and Nomism have largely fallen out of favor. Hypotheticalism has waned because of the implausibility of the idea that anything fundamentally real corresponds to the world-selection function needed for the semantics of the subjunctive conditional. The relative *closeness* of two worlds seems too subjective and anthropocentric to be a metaphysical primitive. Nomism has faded because of the difficulty of bridging the gap between facts about laws and facts about particular patterns of fact. Bridging this gap means attributing an odd sort of *causal power* to the laws themselves. Thus, the two main competitors today are Powerism (or the *powers ontology*) and Neo-Humeanism.

Neo-Humeanism has gradually declined somewhat in popularity as it failed to provide adequate accounts of the directionality of time and causality, of dispositions and powers, of objective probability, and of scientific theory choice and induction (again, see Koons and Pickavance2017). Hence, there has been increasing interest in a Powerist alternative. (Of course, I am not denying that the other three views have their contemporary defenders, nor am I claiming that the issue is a settled one.)

A viable powers ontology must include two additional elements: forms and processes. It is processes that *manifest* powers, and it is forms that *ground* them. Causal powers come in two kinds: active and passive. An active power initiates a process of change (kinesis) in some entity, and a passive power is the potentiality for undergoing such a process.

Powers appear in nature in natural clusters, and these power-clusters are the expression of the presence of Aristotelian *forms* (Inman 2018). Functionally equivalent or interchangeable forms constitute the basis of natural kinds of substances, whether essential or accidental. Without forms as the common ground of these repeatable clusters of powers, we would be left with a large number of massive brute coincidences. The substantial form of water explains why the active and passive powers associated with all instances of water are found so regularly in concert.

Active causal powers initiate ongoing processes of change. Without such processes, it would be impossible to explain how the past influences the future, unless we were to posit immediate action at a temporal distance. Processes of change in turn presuppose the existence of fundamentally enduring entities, the fundamental *participants* in these processes, and these participants must be subject to substantial forms that determine their persistence-conditions and their liabilities to accidental change or motion. Nature’s repertoire of forms determines what kinds of entities are metaphysically fundamental.

109 In contrast, the Neo-Humean ontology requires no fundamental processes or
110 fundamentally enduring entities (with their substantial forms). Instead, what is fun-
111 damental is a framework of spacetime (or spatiotemporal relations), with regions
112 occupied by one or more kinds of qualities or stuffs (the Humean mosaic). Time is
113 metaphysically prior to change, since change is simply a matter of the appearance
114 of different qualities at different times (Russell's At-At theory). Laws of nature are
115 grounded in brute-fact patterns of qualitative succession. On the Mill-Ramsey-Lewis
116 model, a mathematical function counts as a law of nature just in case it is a theorem
117 of the simplest axiomatization of the mosaic's patterns.

118 The two ontologies of causation correspond closely to two philosophies of nature,
119 philosophies that have been in competition since the later Middle Ages. We can call
120 these the *perennial* (or *scholastic*) and the *modern* philosophies. On the perennial
121 philosophy of nature, the task of science is to identify the substantial and acciden-
122 tal forms in nature, from which flow things' active and passive capacities, which
123 manifest themselves (in turn) in the form of activities and processes of change. Math-
124 ematics can be a useful tool in describing these capacities and processes, but science
125 is primarily concerned with discovering the *real definitions* of natural kinds. In addi-
126 tion, the realm of *potentiality* is real and inescapable, even if in some sense dependent
127 on the actual. The reality of potentiality (powers) corresponds to the reality of a kind
128 of teleology: the *natural intentionality* (in George Molnar's phrase) of the real but
129 unmanifested potentialities of nature.

130 The perennial philosophy of nature is pluralistic, in that each kind of form could
131 give rise to a distinct set of active and passive powers. This allowed for the possibility
132 of fundamental entities studied in distinct theoretical domains, including chemistry
133 and biology as well as physics. In fact, I will go even further and argue that the
134 quantum revolution requires us to *demote* the status of microphysical entities, includ-
135 ing particles and fields. We should reverse the usual understanding of *emergence*:
136 it is microphysical phenomena that emerge from the more fundamental domain of
137 chemistry, thermodynamics, and solid-state physics, not vice versa. .

138 On the modern view, science is primarily about discovering fundamental math-
139 ematical relations explain and in some sense *govern* observable phenomena. The
140 task is to find increasingly general and simple formulas, from which all such math-
141 ematical relations can be derived through calculation. The realm of potentiality is
142 unreal or imaginary—merely a result of human thought experiments. Natural reality
143 is exhausted by what actually happens. The modern philosophy of science aspires to
144 be absolutely unitary, discovering a single set of laws that apply to all interactions at
145 all scales. In practice, this translates into the priority of the microscopic realm, since
146 large-scale structures and patterns are nothing more than the sum of their small-scale
147 components.

148 3 The quantum revolution

149 Perhaps the most important and yet often overlooked aspect of the quantum rev-
150 olution is its elevation of physical potentialities to a level of indispensability, as
151 Heisenberg recognized (Heisenberg 1958, p. 41) In modern philosophy of nature,

the realm of potentiality can be treated as something unreal, as a mere mental construction or thought experiment. In quantum mechanics, however, what is merely potentially so has a real impact on what actually happens. This comes out very clearly in Richard Feynman's sum-over-history or path integral formulation of QM. In order to predict what will actually happen, one must compute the probability amplitude corresponding to every possible path of the system from initial to final states.

Since the time of Newton and Leibniz, physicists have had two sets of mathematical techniques for explaining and predicting the motion of bodies. One model, the Newtonian, treats force, mass, and instantaneous acceleration as the metaphysically fundamental properties, relying on vector addition (the quadrilateral of forces) to work out the rate and direction of acceleration for each body. This model takes states and events as the primary reality, with a Russellian at-at theory of motion, and binary forces of attraction and repulsion between simple bodies as the ultimate drivers of physical action. This fits nicely with the microscopic or *microphysicalist* commitments of modern philosophy of science.

The second, analytical or Hamiltonian model, gives primacy instead to energies and processes (trajectories) over instantaneous forces, relying on the conservation of energy and principles of least action, instead of Newton's laws of motion (McDonough2008, McDonough2009). The alternative model begins with the Lagrangian formulation of mechanics, in which whole trajectories are explained via some form of 'least-action' or 'extremal' or 'variational' principle (Yourgrau and Mandelstam1979, pp. 19-23, 164-7; Lindsay and Morgenaw1957, pp. 1336; Lanczos1986, pp. xxvii, 345-6).

In classical mechanics, theorists had a free choice between a Newtonian and a Lagrangian/Hamiltonian model, which each being derivable from the other. With the quantum revolution, the second model becomes obligatory, since the fundamental entities can no longer be imagined to be moving in response to the composition of forces exerted at each moment from determinate distances. Teleology reigns supreme over mechanical forces, as Max Planck noted. (See Planck1936, pp. 119-26; Planck1960; Dusek2001; Thalos2013, pp. 84-6) This provides a second line of support between quantum mechanics and the perennial philosophy.

Finally, quantum mechanics represents the microscopic domain as *incomplete*, in that it ascribes to microscopic entities only a probability of being observed or measured in various states, but it leaves the notions of *observation* or *measurement* without any microscopic definition. This is in sharp contrast to classical mechanics, in which there is no essential reference to anything beyond the locations and momenta of the individual particles. This creates a severe problem for the microphysicalist commitments of modern philosophy of nature, a problem that has come to be known as *the measurement problem*. As we shall see, there is no such problem for the scholastic philosophy of nature and its attendant powers ontology.

4 The fundamentality of composite things 192

The perennial or Aristotelian philosophy of nature has the resources to deny the primacy of mereologically simple entities, whether these are so-called "fundamental" 194

195 particles or field values at spatiotemporal points. In contrast, the modern philoso-
196 phy of nature consciously or unconsciously identifies mereological simplicity with
197 metaphysical fundamentality.

198 I will use the term *substance* to refer to the mereologically composite and meta-
199 physically fundamental entities that are posited by the perennial philosophy. These
200 substances can exist at many different scales: microscopic, mesoscopic, macro-
201 scopic, or even cosmic. They are not, however, among the very smallest things in
202 nature, since they have proper parts than which they are larger. Unlike quantum
203 particles, Aristotelian substances always have definite location and trajectory. Cru-
204 cially, the substances have definite locations even though their quantum parts do not!
205 Substances also have a full complement of determinate, classical properties (cor-
206 responding to superselection sectors in algebraic QM).¹ These classical properties
207 include chemical form, chirality, temperature, entropy, and chemical potential.

208 It is when we look at composite substances (including macroscopic ones) that we
209 see the need for Aristotelian hylomorphism, and not merely the so-called *powers*
210 *ontology* of such recent philosophers as C. B. Martin, George Molnar, or John Heil.
211 For example, Heil holds that the only substances that exist are simple and micro-
212 scopic, corresponding to the *fundamental particles* of contemporary physics (Heil
213 2012, pp. 18–22). Such a non-hylomorphic version of powers ontology is in real ten-
214 sion with the apparent holism of quantum mechanics. In addition, as I will argue in
215 Section 7 below, it fails to provide any solution to the quantum *measurement prob-*
216 *lem*. I will defend a hylomorphic account of substances that is precisely the opposite
217 of Heil's: instead of saying that only particles are substances, I will claim that only
218 non-particles are substances, i.e., that no “fundamental” particles are substances at
219 all.

220 There are several reasons for denying quantum particles the status of metaphysically
221 fundamental substances (see Koons2019 Section 2.4). First of all, when parti-
222 cles are entangled, they lose their individual identities, in much the same way that
223 dollars do when deposited in a bank account. This is reflected in the anti-haecceitistic
224 bias of quantum statistics, in both the Bose-Einstein (for bosons) and Fermi (for
225 fermions) forms (see the chapters in Part I of Castellani1998). Second, in relativistic
226 quantum field theory, even the number of fundamental particles is not an absolute fact
227 but varies according to one's frame of reference (see Fraser2008). Thirdly, particles
228 are wavelike in nature—they are merely excitations in fields, not entities in their own
229 right. In standard (non-Bohmian) versions of quantum mechanics, particles typically
230 lack spatial location and spatiotemporal trajectories. Any particle at any time has a
231 finite probability of being detected anywhere in the universe (Clifton and Halvorson
232 2001). Finally, if particles were substances, then explaining the Einstein-Podolsky-
233 Rosen correlations (which violated Bell's inequality) would require super-luminal
234 causation between widely separated particles—effectively, instantaneous action at
235 great distances.

¹Throughout I will use the term ‘classical’ to refer to properties in the non-trivial center of algebraic models—properties that are mutually commuting, corresponding to superselection rules.

Aristotelian substances, being composite, come in two kinds: homogeneous and heterogenous. The most prominent examples of heterogeneous substances are living organisms. Organisms and other heterogeneous substances (if there are any) have clear spatial boundaries. In the case of homogenous substances, like water or hydrogen gas, the spatial individuation of individual substances would seem to be a matter of convention or speculation. It might be the case that for each natural kind of homogenous substances, there is at each point in time just a single scattered individual, one that exists as long as some of the substance exists somewhere. Local substantial change at the level of homogeneous substances is, however, an empirical matter. Wherever symmetries are broken spontaneously, there is a local substantial change from one substance to another (see Section 6.2).

On the Aristotelian model, parts of substances are metaphysically dependent on the whole. Applying this to quantum mechanics would result in the supposition that the states and locations of quantum particles are wholly grounded in the natures and states of the bodies to which they belong (and not vice versa). We could even go so far as to say that quantum particles have only a *virtual existence* until they come to be manifested in interactions between substances. This accords nicely with the fact that quantum particles lack any individual identity. Quantum statistics (in both the Fermi and Bose-Einstein versions) treats indistinguishable particles as lacking ontological distinctness, in contrast to classical statistics.

Quantum mechanics assigns to particles vectors in a state space, with projections of the vectors on various properties corresponding (via Born's rule) with the probability of our observing the particle's exhibiting the property in question. From the perennial perspective, the quantum representation is a representation of a certain active *power* of the whole substance—a power to manifest a particulate part with certain features in interactions with other substances (in this case, the experimenters and their instruments). The Kochen-Specker theorem of quantum mechanics entails that it is impossible to attribute a full range of determinate properties to these merely virtual entities at all times.

5 Against microphysical reduction 265

The perennial philosophy depends on denying that sciences like chemistry, thermodynamics, and biology are reducible to particle or field physics, since entities that are *reduced* to other entities cannot be metaphysically fundamental, and it is chemical and biological substances and not particles or fields that are fundamental.

Most philosophers of science assume that one theory can be *reduced* to another if the dynamical laws of the former can be derived from those of the latter under certain constraints or conditions (the so-called 'classical' or 'Nagelian' model of reduction). However, this common assumption overlooks the fact that every scientific explanation appeals to *two factors*: dynamical laws and a phase space (including a manifold of possible initial conditions). Consequently, every scientific theory comprises two elements: a set of dynamical laws and a space of possible initial conditions. The structure of this space implicitly encodes crucial nomological information.

278 In order to secure a metaphysical conclusion about dependency between the
279 domains of two theories, it is not enough to derive the dynamical laws of one theory
280 from the dynamical laws of the other, supposedly more fundamental theory. We must
281 also prove that the structure of the phase space and of the manifold of possible initial
282 conditions of the supposedly reducing theory is not itself grounded in the structure
283 or laws of the reduced theory.

284 Suppose, for example, that we have two theories, T_1 and T_2 . Theory T_1 consists in
285 a set of dynamical laws D_1 together with a phase space S_1 , and T_2 similarly consists
286 of laws D_2 and space S_2 . Let's suppose that we have a Nagelian reduction of T_1 to
287 T_2 : a translation $*$ from the vocabulary of T_1 into T_2 such that D_2 entails $(D_1)^*$ with
288 respect to space S_2 , but $(D_1)^*$ does not entail D_2 with respect to S_2 : that is, the set
289 of trajectories (the flow) through S_2 that are logically consistent with D_2 is a proper
290 subset of the set of trajectories through S_2 that are consistent with $(D_1)^*$.

291 Would this narrow or Nagelian "reduction" give us grounds for taking the entities
292 and properties of T_1 to be wholly *grounded* in those of T_2 ? Not necessarily: we have
293 to take into account the role of the phase spaces S_1 and S_2 . Suppose, for example,
294 that the structure of S_2 (the supposedly reducing theory) is metaphysically grounded
295 in the structure of S_1 : it is facts about the natures of the supposedly reduced theory T_1
296 that explains the structure of the space of possibilities used to construct explanations
297 in terms of theory T_2 . It may be, for example, that the structure of S_1 is "tighter" or
298 more restrictive than the structure of S_2 (under any metaphysically sound translation
299 between the two), and this tighter structure might be inexplicable in terms of D_2 ,
300 theory T_2 's dynamical laws. Space S_1 could have additional *structure*, in the form
301 of new, irreducible properties. In addition, there might be no natural restriction on
302 space S_2 that would close the modal gap between S_1 and S_2 . On these hypotheses, the
303 Nagelian reduction of the dynamical laws of T_1 to T_2 would carry no metaphysical
304 implications.

305 It was easy to overlook this fact, so long as we took for granted the ungrounded
306 and even universal nature of the microscopic or microphysical phase space. In classical
307 mechanics, the space of possible boundary conditions consists in a space each of
308 whose "points" consists in the assignment (with respect to some instant of time) of
309 a specific location, orientation, and velocity to each of a class of micro-particles. As
310 long as we could take for granted that this spatial locatedness and interrelatedness of
311 particles is not metaphysically grounded in any further facts (including macroscopic
312 facts), reduction of macroscopic laws to microscopic dynamical laws was sufficient
313 for asserting the complete grounding of the macroscopic in the microscopic, and
314 therefore for asserting the ungroundedness (fundamentality) of the microphysical
315 domain. However, this ungroundedness of the spatial locations of microscopic parti-
316 cles is precisely what the quantum revolution has called into question. As I will argue
317 in Sections 6 and 7 below, the phase space of macroscopic objects involves classical
318 properties that cannot be derived from the non-commuting, quantal properties of
319 pure quantum mechanics. The introduction of the thermodynamic or continuum limit
320 introduces new mathematical structure to the phase space of thermodynamics, rendering
321 the metaphysical reduction of thermodynamics to particle physics impossible,

even though the dynamic law governing thermodynamics (the Schrödinger equation) is wholly derived from particle physics. 322
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6 Thermochemical powers and potentialities 324

From² the 1950's onward, quantum theory moved from the pioneer period to that of generalized quantum mechanics. Generalized QM moved away from the Hilbert-space representation of pure quantum systems to an algebra, in which both quantum and classical observables could be combined in a single formal representation. The algebras of generalized QM can have non-trivial *cores*, consisting of the classical properties that commute with every other property, representing exceptions to the mutual complementarity of the quantum variables. In practice, this means representing the classical properties of complex systems (like molecules or experimental instruments) as ontologically fundamental, on par with the quantum properties of the smallest particles. 325
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In addition, by moving to the “thermodynamic” or continuum limit, which involves treating a system with apparently finitely many parameters or degrees of freedom as though there were infinitely many such degrees, algebraic QM enabled theorists to introduce superselection rules, which could be used to distinguish the different phases of matter that can co-exist under the same conditions (such as gas, liquid, solid, ferromagnetized, superconducting). I will argue in the following subsections that the use of the continuum limit can best be interpreted as representing an ontological difference between two irreducibly macroscopic conditions, providing strong evidence against reduction. 335
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6.1 The continuum limit: a mark of ontological fundamentality 344

In applied physics, it is common to take some parameter to infinity: that is, to replace the original model having some finite parameter with a new model in which that parameter takes the value of infinity. For example, in the so-called “thermodynamic” limit, a system containing n molecules and a fixed volume V is replaced by one in which both the number of molecules and the volume go to infinity, while keeping the density n/V constant. As Compagner explains (Compagner 1989), this thermodynamic limit is mathematically equivalent to the *continuum limit*: keeping the volume constant and letting the number of molecules go to infinity, while the size of each molecule shrinks to zero. In many applications, such as the understanding of capillary action or the formation of droplets, the continuum limit is the right way to conceptualize the problem, since infinite volumes have no external surfaces and cannot interact with their containers. 345
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As Hans Primas has pointed out (Primas 1983), there are three reasons for taking infinite limits in physics: for mathematical convenience, in order to isolate some fac- 357
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²In this section, I build on my own work in (Koons 2018b) and (Koons 2019). See also the recent work by William M. R. Simpson: (Simpson 2020, Chapter 7).

359 tors from others, and in order to introduce new structure into the representation. The
360 continuum limit in generalized quantum mechanics is an example of the third reason.
361 In 1931, John von Neumann and Marshall Stone proved that finite systems admit
362 of only one irreducible Hilbert-space representation (Neumann 1931). Infinite sys-
363 tems, in contrast, admit of infinitely many inequivalent Hilbert-space representations.
364 This apparent embarrassment of riches in the infinite case turns out to be crucial for
365 the representation of phase transitions, entropy, and thermodynamic phenomena. As
366 Geoffrey Sewell explains:

367 For infinite systems, the algebraic picture is richer than that provided by any
368 irreducible representation of observables. . . Furthermore, the wealth of inequiv-
369 alent representations of the observables permits a natural classification of the
370 states in both microscopic and macroscopic terms. To be specific, the vec-
371 tors in a [single Hilbert] representation space correspond to states that are
372 macroscopically equivalent but microscopically different, while those carried
373 by different [inequivalent] representations are macroscopically distinct. Hence,
374 the macrostate corresponds to a representation and the microstate to a vector in
375 the representation space. (Sewell 2002, pp. 4-5)

376 Thus, at the thermodynamic limit, algebraic quantum mechanics gives us exactly
377 what we need: a principled distinction between quantal and classical (non-quantal)
378 properties. In addition, the non-quantal properties do not supervene on the quantal
379 properties of a system, since the latter always consists of a *finite* number of facts,
380 while the thermodynamic limit requires an infinite number of virtual sub-systems.
381 The classical features are real and irreducible to the quantum particle basis. As I will
382 argue in Section 7, this is exactly what is needed to resolve the quantum measurement
383 problem.

384 Franco Strocchi (Strocchi 1985) has shown that the continuum limit is needed to
385 explain any spontaneous symmetry breaking in quantum-mechanical terms. In classi-
386 cal mechanics, symmetry breaking could always be explained by small perturbations
387 with non-linear consequences. These small perturbations or prior asymmetries can be
388 ignored for the sake of convenient, approximate representations. In quantum mechan-
389 ics, this simply does not work. Strocchi points out that in many cases “it is impossible
390 to reduce symmetry breaking effects to asymmetric terms in the Hamiltonian.”
391 (Strocchi 1985[p. 117]) The dynamics have to be defined in terms of a symmetric
392 Hamiltonian. Consequently, we need true emergence of asymmetry, not simply the
393 apparent emergence that results from suppressing slight asymmetries in the prior
394 situation (as in classical mechanics). This is possible only for infinite quantum
395 mechanical systems. Any finite system retains any symmetry that it possesses.

396 In addition to symmetry breaking, infinite algebraic models are also crucial to the
397 representation of irreversibility, which, in turn, is essential to thermodynamics (as
398 noted by Woolley Woolley 1988, p. 56). This reflects work by Ilya Prigogine and his
399 collaborators, who demonstrated that molecular motions any finite quantum system
400 are always perfectly reversible. This is not the case for infinite systems, which can
401 show irreversible behavior and thus can validate the Second Law of Thermodynamics
402 as a fundamental law of nature.

6.2 Thermodynamics and phase transitions: infinite algebraic models

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The infinite algebraic models of generalized QM provide, for the first time, the possibility of rigorous and non-arbitrary definitions of the basic thermodynamic properties of entropy, temperature, and chemical potential see (Sewell 2002). Contrary to what many philosophers believe, science does not suppose that temperature is the mean kinetic energy of molecules! (Vemulapalli and Byerly 1999, pp. 28-30) See also (Primas 1983, pp. 312-3).

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If the system is not at equilibrium, temperature is not well-defined, though the mean kinetic energy is. . . . Temperature is a characteristic of equilibrium distribution and not of either individual molecules or their kinetic energy. When there is no equilibrium between different kinds of motion (translations, rotations, and vibrations), as in the case of molecular beams, temperature is an artificial construct. (Vemulapalli and Byerly 1999, pp. 31-2)

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Since thermal equilibrium is not defined at the level of statistical mechanics, temperature is not a mechanical property but, rather, emerges as a novel property at the level of thermodynamics. (Bishop and Atmanspacher 2006, p. 1769)

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If temperature could be defined as mean kinetic energy, then temperature would always be defined for any collection of molecules, since the kinetic energy of each molecule is always well-defined. In fact, many physical bodies have no well-defined temperature, as Vemulapalli and Byerly point out in the above quotation. Temperature emerges only once a thermodynamic equilibrium has been established between different modes of kinetic energy. Thus, without the thermodynamic limit as a faithful representation of real systems, we would have to dismiss all talk of 'temperature' as merely a useful fiction.

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In addition, *phase transitions*, such as those between the solid, liquid, gas states, and between conditions before and after the onset of coherent ferromagnetism or superconductivity in metals, require the use of infinite models (models involving the continuum limit): see (Liu 1999), (Ruetsche 2006), and (Bangu 2009). Phase transitions are an important case of spontaneous symmetry breaking (Sewell 1986, p. 19).

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6.3 Molecular structure

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Generalized quantum mechanics attributes both classical and quantum properties to objects. The modern quantum theory of molecular structure is a classic example. The structure of a molecule, that which distinguishes one isomer from another, including right-handed chiral molecules from left-handed ones, depends entirely on the classical properties of precise location applied to atomic nuclei. As Hans Primas put it, "Every chemical and molecular-biological system is characterized by the fact that the very same object simultaneously involves both quantal and classical properties in an essential way. A paradigmatic example is a biomolecule with its molecular sta-

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443 bility, its photochemical properties, its primary, secondary, and tertiary structure.”
444 (Primas1983, p. 16) . The quantal properties of a system correspond to the wavefunc-
445 tions associated with each of its constituent particles. These wavefunctions play a
446 crucial role in explaining the behavior of bonding or valence electrons in molecules,
447 as well as such phenomena as super-conductivity (Cooper pairs of electrons) and
448 super-fluidity.

449 **7 Powers and the measurement problem**

450 Pioneer quantum mechanics is pure quantum mechanics, in the sense that all
451 (or nearly all) observables are quantum observables—mutually complementary (in
452 Bohr’s sense), satisfying the superposition principle. A classical observable is a prop-
453 erty that commutes with all other properties, meaning that it can be conjoined, in a
454 physically meaningful way, with any other observable. An entity can have a deter-
455 minate value of a classical observable at all time, while it is impossible to have
456 determinate values for two, mutually non-commuting quantum observables. As an
457 expression of this pioneer viewpoint, John von Neumann laid down the *irreducibility*
458 *postulate* (Neumann 1931): no two observables are commutative.

459 Irreducibility gives rise inevitably to the so-called “measurement problem”: exper-
460 iments invariably take place in a context defined in terms of classical observables,
461 like location and temperature. If the theory includes no classical observables, then
462 there is an unbridgeable conceptual gap between the world of theory and the world of
463 the experimenter. The different responses to the measurement problem produced the
464 different “interpretations” of the formalisms of Pioneer Quantum Mechanics. Here
465 are the five most common and well-defended interpretations:

- 466 1. The Copenhagen interpretation or family of interpretations, comprising a vari-
467 ety of pragmatic, operationalist, perspectivalist, and anti-realist interpretations,
468 including that of Niels Bohr. Quantum states are defined in terms of experimental
469 results and have no independent existence.
- 470 2. Dualist interpretations: Eugene Wigner, John von Neumann. Human conscious-
471 ness causes a “collapse of the wave packet”: a discrete transition from a
472 superposed quantum state into a state in which the system possesses some
473 definite value of the appropriate classical property (position, momentum, etc).
474 This involves positing two distinct dynamics in the world—one occurring
475 autonomously, the other existing in response to interactions with consciousness.
- 476 3. David Bohm’s interpretation (Bohm 1951), building on Louis de Broglie’s 1925
477 pilot wave account. The pure quantum world exists with a unified, uninterrupted
478 dynamics. The universe consists of point particles with definite locations at all
479 times, guided by the wave function, and forming a single, indivisible and non-
480 localizable dynamical system.
- 481 4. Hugh Everett’s (1957) “relative state” or “many worlds” interpretation, devel-
482 oped by Bryce De Witt, R. Neill Graham, David Deutsch, and David Wallace
483 (Wallace 2008). The classical world of experiments is merely an appearance,

a product of the limited perspective of human and other organisms. When performing experiments involving interaction with systems in superposed quantum states, the observer splits into multiple versions, one corresponding to each possible state. Each split state involves no awareness or memory of states experienced in parallel branches.

5. Objective collapse theories, such as GRW (Ghirardi et al. 1985). These interpretations are like the dualist versions, except that the collapse of the wave packet is triggered by certain physical events and not by consciousness. At this point, these theories go beyond interpretation, postulating a new, so-far merely speculative collapse-triggering mechanism. At this point, there is no specific theory and no empirical confirmation. In addition, objective collapse theories require still further ontological interpretation, such as John Bell's "flash ontology" (Bell 1987) or the matter density model.

Hylomorphism with its power ontology can be offered as a sixth interpretation, an interpretation inspired by some remarks of Heisenberg (Heisenberg 1958), and defended by Nancy Cartwright (Cartwright 1999) and Hans Primas. Interaction between the quantum powers of one substance and the substances making up the experimenters and their instruments precipitates an objective collapse of the quantum object's wavefunction, as a result of the joint exercise of the relevant causal powers of the object and the classical instruments,³ and not because of the involvement of human consciousness.

How is this a solution to the measurement problem? Why haven't I merely restated the problem by referring to 'observers' and their 'classical instruments'? My answer is this: according to hylomorphism, observers and their instruments are substances (or made of substances), and substances are not composed of quantum particles. The states of substances are not reducible to the quantum states of their particles. Thus, there is no inconsistency in supposing that substances have properties ('classical') that are exempt from superposition and that, therefore, always constitute definite outcomes. I will explain how this works in more detail in Section 7.2 below, following the work of Hans Primas.

Do we need perennial philosophy and not just some version of contemporary powers ontology? Yes, because if we try to solve the measurement problem with powers alone, we will have to attribute those powers to quantum particles and only to quantum particles. This would include both active and passive powers. Solving the measurement problem requires observers and their instruments to have *non-quantal passive powers*, through which they can register definite results and not merely enter into an extended superpositions. As I have argued above, Aristotelian substances have the capacity to bear irreducible chemical and thermodynamic properties (as represented in the non-trivial centers of infinite algebraic models). Quantum particles do not have that capacity: they are fully characterized by vectors in a single Hilbert space in a finite algebra with only a trivial center and no superselection sectors.

³This joint exercise is an instance of what is known in the causal powers literature as *mutual manifestation*: see (Heil 2003) and (Mumford and Anjum 2011).

525 7.1 Epistemological constraints on a solution to the measurement problem

526 To solve the measurement problem, it is not enough for an interpretation of quan-
527 tum mechanics to merely *save the phenomena*, in the sense of merely explaining
528 how it is possible for us to experience the appearance of a macroscopic world (with
529 objects instantiating mutually commuting, *classical* observables like actual position).
530 We must distinguish between *explaining* and *explaining away*. A credible scientific
531 theory must explain most of our apparent data, in the sense of both treating it as
532 objectively known fact and providing a satisfactory causal account of its genesis. A
533 scientific theory that *explains* the data by entailing that it is all a mere appearance,
534 without objective reality, destroys its own empirical foundations.

535 More specifically, here are some epistemological constraints that must be satisfied
536 (see Simpson Simpson2020, Chapter 8; Simpson2019):

537 **E1. Perception.** The theory must endorse the fact that our sensory perception of
538 physical events and objects is mostly reliable.

539 **E2. Memory.** The theory must endorse the fact that our memory of past observa-
540 tions is mostly reliable.

541 **E3. Induction.** The theory must endorse the fact that the physical events and facts
542 that we observe (currently and in memory) are an inductively reliable sample of
543 the whole.

544 As we shall see, each of the new interpretations of QM fails one or more of these
545 tests, in contrast to the power ontology of hylomorphism.

546 The non-locality of quantum mechanics, as exemplified by Bell's theorem, threat-
547 ens condition E1. If we embrace a Neo-Humean account of causation, the immediate
548 consequence is that causation in the quantum domain is radically non-local. By *rad-*
549 *ically non-local*, I mean that the intensity of the influence of distant bodies does not
550 decrease as distance increases. Very remote objects (if entangled with something in
551 our neighborhood) can have effects every bit as significant as other objects in that
552 same neighborhood. In principle, at least, this raises questions about the reliability of
553 our sensory perception of our immediate environment, since our brains or our sense
554 organs might be entangled with distant objects in a way that makes them unreliable
555 as indicators of local conditions.

556 Hylomorphists can secure the justifiability of reliance on perception by posit-
557 ing receptive causal powers that, when not interfered with by abnormal conditions
558 (whether internal or external), actualize themselves in the form of veridical impres-
559 sions of one's environment. Since Neo-Humeans lose such a robust Aristotelian
560 theory of causal powers, with its distinction between normal and abnormal condi-
561 tions, they are left with a situation in which the fallibility of the sensory process
562 makes it unreasonable to treat any sensory impression as knowledge-conferring.

563 7.2 The neo-copenhagen (hylomorphic) programme

564 The old Copenhagen view of Niels Bohr suffered from being too narrowly dualistic,
565 distinguishing the classical world from the quantum world. In contrast, the hylomor-
566 phic interpretation embraces a salutary kind of ontological pluralism, recognizing

that the non-quantum or supra-quantum world is itself a “dappled” world (as Nancy Cartwright puts it), dividing naturally into multiple domains at multiple scales. This fits the actual practice of scientists well, who are in practice ontological pluralists, as Cartwright has documented.

The measurement problem arises from the formulation of quantum mechanics as a theory about the probabilities of certain measurement results. The quantum wavefunction evolves in a deterministic manner, by the unitary dynamics of Schrödinger’s equation. In order to test the theory, some observable results must be deduced from the theory. It is Born’s rule that enables us to move from some parameter value in the wavefunction (the wave amplitude) to something testable: namely, certain probabilities about the result of measuring one or other classical parameter (such as position or momentum). This early model (as developed by Bohr and Heisenberg) assumed that we could continue to use classical language in describing the experimental setup and the measurement devices. Critics have argued that this involves an implicit inconsistency, since physicists assume that these classical instruments are wholly composed of quantum systems and so should be, in principle, describable in purely quantum and not classical terms.

This charge of inconsistency falls flat when lodged against the hylomorphic version of the Copenhagen programme. Observers and their instruments are not reducible to their quantum constituents—instead, quantum particles have only virtual existence, corresponding to certain powers of thermochemical substances. Theoretically, this depends (as I showed in the last section) on the use of algebraic formulations of quantum mechanics with infinite models (at the continuum limit). The additional structure afforded by such models represents the irreducible fundamentality of these substances.

Bohr’s interpretation required that reality be divided into two disjoint realms, the classical and the quantum, with a measurement involving any setup in which a quantum system is made to act upon a classical observer or instrument. This foundered on the fact that some systems, like supercooled fluids or quantum computer chips, bridge the gap between the two realms. We cannot consistently describe all macroscopic objects in purely classical terms, as Bohr’s program seems to require, since it is interaction with the classically described realm of measurement devices that collapses the wavefunction in Bohr’s model. In contrast, on the Primas model, we could postulate that the wave packet associated with a quantal property has “collapsed” whenever it becomes correlated with a fundamental *classical property* of a disjoint system. Even though entities cannot be neatly divided into two disjoint domains, this is not true of physical properties. Substances have *both* classical properties *and* (by virtue of their virtual parts) quantal properties. Infinite algebraic models represent quantal properties as vectors in individual spaces and classical properties as disjoint spaces or superselection sectors.

Primas demonstrates (Primas 1990) that interaction with the classical properties of entities in the environment will drive quantal vectors to eigenstates with a high probability in a short period of time. The Primas solution is, consequently, one of continuous rather than discrete collapse (unlike, for example, most versions of the GRW model of objective collapse). The Primas model can be incorporated into a

612 powers ontology, by attributing to substances the power to collapse the wavefunctions
613 associated with quantum parts of other substances.

614 Bell characterized the measurement succinctly in this way: either the Schrödinger
615 equation isn't right, or it isn't everything. Most solutions to the problem fall squarely
616 into one side or the other: the Copenhagen interpretation and the many-worlds inter-
617 pretation insist that the equation isn't everything, while the GRW and other objective
618 collapse theories suppose that it isn't right. On which side does hylomorphism stand?
619 I've described it as a neo-Copenhagen view, while Primas offers a model of objective
620 collapse.⁴

621 Of course, Bell's alternatives are not exclusive. In fact, the Schrödinger equation
622 is neither everything nor right. It is right insofar as it describes the evolution of the
623 quantal aspects of a substance sans interaction with other substances. However, this is
624 not everything, since thermal substances also possess determinate, non-quantal prop-
625 erties. And it is incorrect, even as a description of those quantal aspects, whenever
626 the quantum potentialities are actualized through interaction with other substances.
627 At that point, a form of objective collapse takes place, in a way described by Primas's
628 model.

629 **7.3 The everettian programme**

630 There⁵ are three defects to the Everett (relative-state or branching world) programme,
631 each of which hylomorphism avoids. First, hylomorphists can give a straightfor-
632 ward, intuitive, and natural account of the *probabilities* associated with the quantum
633 wavefunction: the square of the wave's amplitude associated with some precise state
634 represents the probability that the quantum particle will interact in a corresponding
635 way with some classical measurement instrument. So, for example, if we use a pho-
636 tographic plate to register the location of a photon, then the quantum probability
637 associated with a particular location will give us the probability that the photon will
638 interact with the plate at that location. In contrast, the Everett interpretation requires
639 that we radically modify our naïve conception of probability, assigning fractional
640 probabilities to various states, even though it is certain that each of the states will
641 in fact be realized (although on different "branches" of the world). See (Kent 2010;
642 Price 2010). I have argued that the sophisticated, neo-pragmatist solution to this
643 problem developed by David Wallace and other "Oxford Everettians" fails, because
644 it overlooks the possibility of a rational agent's utility depending on inter-branch
645 comparisons (Koons 2018a).

646 The second drawback to the Everett interpretation is that it, like the Bohm inter-
647 pretation, renders our classical interactions with the quantum world illusory. There
648 are, on the Everett interpretation, no *inter*-actions at all. The evolution of the world
649 is simply the autonomous unfolding of a single object, the universe, according to
650 a global Schrödinger equation. Entities like you and I and our experimental instru-
651 ments are merely simulated by aspects of this function, as a kind of "virtual reality".

⁴Thanks to an anonymous reviewer for pressing this question.

⁵This section builds on my work in (Koons 2018a). See also (Simpson 2020, Chapter 8)

(See Albert²⁰¹⁵, Halliwell²⁰¹⁰, Maudlin²⁰¹⁰) The world has all the causal oomph there is, leaving nothing over for mere parts of the world to exercise. This means that the Everett interpretation must lose all of the epistemological advantages that a causal-powers account of scientific experimentation can provide.

In effect, the Everett interpretation (in its modern, Oxford-school form, as developed by David Wallace ²⁰⁰⁸ and his collaborators) almost perfectly duplicates Plato's allegory of the cave from *Republic* Book VI: we are forced to watch the mere shadows (the classical observables) cast by the quantum wavefunction, which lies always outside our field of vision. In fact, we are in an even worse predicament than the prisoners in the cave: since we (the observers) *are also mere shadows* on the cave wall. The classical world consists of mere shadows shadow-observing other shadows, with no real entities to whom the appearances can appear. In contrast, the hylomorphic interpretation is fully compatible with attributing real and fundamental causal powers both to the classical and to purely quantum objects.

Is this really fair to the Oxford Everettians?⁶ They could plausibly claim that, on their view, the manifest or classical world is *real* although not *fundamentally* so. It seems unfair to compare the manifest world on their account with virtual reality or with the shadows on Plato's cave. The manifest world is a *real pattern* (to use Daniel Dennett's phrase, Dennett¹⁹⁹¹), one that is functionally realized by the underlying quantum reality. As we shall see (when we turn to my third objection), there are many patterns to be found in the quantum wavefunction. Every logically consistent story with the right cardinality is functionally realized by the quantum world. Therefore, the classical world of experimenters and their instruments is no *more* real than any fiction.

Thirdly and finally, the Everett interpretation leads to global skepticism via both Putnam's paradox (Putnam¹⁹⁸⁰, Lewis¹⁹⁸⁴) and Goodman's grue/bleen paradox (Goodman¹⁹⁵⁴, Lewis¹⁹⁸³), as I have argued elsewhere (Koons ^{2018a}). Putnam's paradox is an argument that purports to show that our words and concepts cannot pick out determinate properties, since the finite class of actual attributions of those words and concepts radically under-determines their extension with respect to not-yet-encountered instances. The standard response to this paradox is to appeal to the relative naturalness of properties whose relevant sub-extension matches our actual use: our words or concepts pick out that most natural property (if there is one) whose extension and anti-extension best fits our actual use of the word or concept in particular affirmations and denials. However, the Everett interpretation is committed to the radical non-naturalness of all the properties that putatively apply to entities in our familiar spacetime world. Hence, our concepts and words can be matched to the truly natural properties (those instantiated by the quantum wavefunction) in an infinite number of equally natural ways. (This is a generalization of an argument by Bradley Monton against wavefunction realism: (Monton ²⁰⁰²) and (Monton ²⁰⁰⁶).)

Suppose that we have two Everettian models of the universe, M_1 and M_2 , with the same cardinality, where each model assigns a Hilbert vector in the same space

⁶Thanks to an anonymous referee for this objection.

694 to each moment of time. (I'll assume that the spacetimes of the two models are iso-
695 morphic.) Let's suppose that M_1 represents the underlying microphysical reality of
696 our actual universe and M_2 that of an alternative, fantastical universe (like Tolkien's
697 Middle-Earth). Let's also suppose that the unitary time-operators and the Schrödinger
698 equations for the two models are both linear and deterministic, although they may be
699 otherwise quite different. Then there are models M_1^* and M_2^* and homomorphisms
700 H_1 and H_2 from M_1^* to M_1 and M_2^* to M_2 (respectively), where M_1^* consists of
701 the representation of an approximately classical, macroscopic, $3 + 1$ -dimensional
702 world that corresponds to the common-sense history of our phenomenal world, and
703 M_2^* a similar representation of the fantastical history (with terms in the Hamiltonian
704 representing the effects of wizardry, for example).

705 There will be a bijective function J (given the linearity and determinism of the
706 dynamics of quantum mechanics) between the vectors of M_1 and M_2 , which pre-
707 serves the underlying dynamics (in the sense that a dynamically possible trajectory
708 in M_1 will be mapped onto a dynamically possible trajectory in M_2). Mapping J will
709 then preserve the truth-values of the microscopic counterfactual conditionals of the
710 two models, so long as the antecedents of the conditionals specify complete states
711 of the universe. In addition, the composition of H_2 and J will be a homomorphism
712 from M_2^* into M_1 . Let's assume, further, that the closeness of two world-states (from
713 a macroscopic perspective), for the purposes of evaluating counterfactual condition-
714 als relevant to M_1^* and M_2^* , is indifferent to the underlying microscopic models. If
715 so, we can adopt a measure of counterfactual closeness on the states of M_1 that per-
716 fectly preserves, under H_2 composed with J , all of the phenomenal and macroscopic
717 counterfactuals true in M_2^* (see Lewis2001). Hence, our actual universe will contain
718 implicitly a representation of the fantastical history M_2^* , in *exactly the same sense* in
719 which it contains a representation of our common-sense history M_1^* .

720 If the only conditions on the extraction of a *phenomenal* or *quasiclassical* world
721 from the wavefunction are mathematical (i.e., the existence of some isomorphism and
722 some measure of closeness that jointly preserve dynamics and the truth-value of condi-
723 tionals), then any imaginable world can be extracted from any wavefunction. The
724 world of Greek mythology, *The Matrix*, *The Lord of the Rings*, or *Alice and Wonder-*
725 *land* would be every bit as real as the world represented in our science and history
726 textbooks. There would be minds experiencing an infinite variety of phenomena,
727 the vast majority of which would have no correspondence whatsoever to the classi-
728 cal physics of Newton and Maxwell. Inhabitants of these non-classical phenomenal
729 worlds would have no hope of ever discovering the fundamental laws of physics.

730 The only way to block these conclusions is to claim that the homomorphism H_1
731 preserves the naturalness of macro properties, the real causal connections between
732 macroscopic things, or the real closeness between states of the world in a way that
733 the composition of H_2 with J does not. However, on the Everett view, there are *no*
734 *natural properties* and *no real connections* in our phenomenal world, and the laws of
735 quantum mechanics do not dictate which pairs of states are really closer than others
736 for the purposes of evaluating macroscopic conditionals, and hence there is no basis
737 for preferring one homomorphism over another.

738 Reflection on these facts would, in turn, provide us with an effective defeater
739 of our own scientific beliefs, since the vast majority of minds would be radically

deceived about the deep nature of the world they (and we) really inhabit, and we would have no non-circular grounds for believing that we inhabit one of the few epistemically “lucky” phenomenal worlds. 740
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Everettians could respond by insisting that the only real branches (the only ones inhabited by really conscious beings) are those that approximate the dynamics of classical physics. In fact, many recent Everettians have implicitly made just such a stipulation: (Albert 1996, pp. 280-1; Gell-Mann 1996; Lewis 2004, p. 726). However, this would be a purely ad hoc move, with no plausible rationale. It would outrageously parochial and anthropocentric, given our own entirely derivative status in the Everettian universe.⁷ 743
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The problem of multiple domains also puts at risk the rationality of induction as a guide to the future. Even assuming that our own domain has been approximately classical up to this point in time, there are many, equally natural extensions of that domain into the future, most of which invalidate our inductive expectations. This involves the application of Nelson Goodman’s grue/bleen paradox to the problem of extracting domains from the wavefunction. In Goodman’s thought-experiment, we are to imagine a possible future in which emeralds continue to be grue, rather than green, after the year 2020 (where ‘grue’ is defined as ‘green if discovered before 2020, and blue otherwise’). Goodman argues that our inductive experience with emeralds before 2020 gives us equally good reason to believe the hypotheses that all emeralds are grue and that all are grue. 750
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When transferred to the Everettian scenario, Alberto Rimini (Rimini et al. 1979) has shown that we can find actual domains in which objects shift in their behavior with respect to a standard set of observables but remain the same with respect to some gerrymandered, “gruesome” observables. Each consistent branch in the Everett multiverse has multiple extensions into the future corresponding to different observable-operators. Some of these extensions are intuitively *unnatural*, in the sense of treating grue-like objects as qualitatively the “same,” before and after the crucial transition. These alternative future branches of our domain are equally natural from the perspective of the underlying quantum wavefunction. Hence, the Everettian has no grounds for privileging what we would deem the more natural branch, since true naturalness must be wholly grounded in what is metaphysically fundamental. 761
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The link between naturalness and fundamentality If instantiations of F and G are wholly grounded in instantiations of (respectively) fundamental properties F' and G' , then if F is more natural than G , so too F' must be more natural than G' . 772
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Goodman’s grue/bleen paradox can be taken as a special case of the Putnam paradox: one in which it is indeterminate how to extend our empirically well-confirmed hypotheses into the future, across an arbitrarily chosen boundary. 775
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These grue/bleen-like paradoxes pose a dilemma for the Everettians. If they suppose that there is no natural mapping from our concepts to features of the real wavefunction, then they have to embrace a radical indeterminacy of interpretation 778
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⁷Schlosshauer (Schlosshauer 2005, p. 1299) points out, “It has become clear that most consistent histories are in fact flagrantly nonclassical.”

781 that deprives nearly all of our assertions and beliefs of determinate truth-value. If,
782 alternatively, they suppose that there is some brute semantic matter of fact about
783 the correspondences, then they have to embrace a scenario in which our inductive
784 practices are radically unreliable, since each empirical generalization will be falsi-
785 fied in many such interpretations, and the Everettians have no grounds for supposing
786 that the one “correct” interpretation is one that verifies the majority of our inductive
787 inferences, bringing the Everett interpretation into conflict with E3.

788 But what about Dennett’s *real patterns*? (Dennett 1991) Couldn’t we insist that
789 our classical world is a real pattern, and that all of these other fictions are merely
790 unreal? What makes a pattern *real*, in Dennett’s account? Dennett says that a pat-
791 tern is real when it is “readily discernible” or “recognizable” (Dennett1991, p. 33).
792 The reality of a pattern depends on “perceivers’ capacities to discern patterns” (Den-
793 nett1991, p. 34). We create real patterns by bringing our pattern-making perspectives
794 to the buzzing blooming confusion of data. (Dennett1991, p. 36) Finding real pat-
795 terns enables us to engage in efficient and reliable prediction. (Dennett1991, p. 42)
796 There is one central problem with all of this: we, with our pattern-recognizing and
797 pattern-making capacities, are also *part of* the very manifest world that we are trying
798 to distinguish from merely fictional patterns. Dennett’s account is either viciously
799 circular or tacitly dualistic, assuming that we exist as real observers outside of the
800 quantum reality whose patterns we recognize. Hylomorphism enables us to avoid
801 such implausible mind/body dualism.

802 **7.4 The bohmian programme**

803 Like the Bohm view, the hylomorphic interpretation assumes a broadly realist stance
804 toward the classical world. Bohm takes classical objects to be composed of particles
805 really located (for the most part) in the regions of space that they appear to occupy
806 in our experience. A deterministic version of Bohm’s theory would seem to offer
807 Neo-Humeans and microphysicalists their best chance at surviving the quantum rev-
808 olution. Each particle in Bohm’s theory has a definite location at each time, and these
809 locational states are indeed fully separable. Each particle has its own unique identity,
810 blocking any quantum fusion.

811 However, there are real concerns about whether Bohm’s theory can underwrite
812 the reliability of our perception of the positional states of our measuring devices.
813 Our subjective impressions would seem to depend on the contemporaneous states
814 of our brains, not the positions of particles in our measuring devices (or even our
815 sense organs, like the retina). Bohm’s theory is certainly capable of generating false
816 sense impressions and false memories about particle positions, since particles do not
817 influence each other’s positions, but are always guided by the cosmic wavefunction.

818 Here’s the form of the argument:

- 819 1. To be empirically adequate, Bohm’s theory must give an account, not just of the
820 “pointer settings” of measuring instruments, but also of our perceptions of those
821 settings (as Bohm himself admitted, Bohm1951, p. 583).

2. There is good reason to think that mental states aren't determined by particle positions within the brain alone. We must include all of the functional features of the brain. 822
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3. But this requires that the basis of mental states includes the state of the cosmic wavefunction, which leads to the radical non-locality of the relevant brain state. 825
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4. In the absence of pervasive and stable decoherence linking brain states and sensible objects, functional states of those states in relation to the brain do not fix particle positions (in either the object or the brain): two pairs of brain-object relational states can be functionally indistinguishable, even though they involve radically different particle positions and trajectories. Therefore, in the absence of effective decoherence, one and the same system (e.g., the person's brain plus his sense organs) cannot be reliable both at tracking functional states and at tracking particle positions. 827
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5. Non-local quantum effects threaten to destroy any reliable correlation between the functional states of the environment and local particle positions and therefore to destroy any correlation between brain states and particle positions. 835
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6. This could be avoided only if we had good grounds for assuming that environmental interaction secured (through decoherence) the effective classicality of the brain-environment interaction, but that is very much in dispute. In addition, Bohm's theory raises special technical problems for the widespread application of decoherence (see Schlosshauer2005, p. 1297-8 and Simpson2019). 838
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7. Evolution would explain our ability to track reliably the relevant *functional aspects* of our environment, not our ability to track particle positions. Evolution cares about whether we can survive and reproduce—it is completely indifferent to whether we can reliably track particle positions. 843
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Brown and Wallace explain why the perceptual state must be fixed by the functional state of the brain, not just by its configuration of particles (premise 2): 847
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Observables in the context of Bell's remark are defined relative to sentient observers, and it is a tenet of the de Broglie-Bohm picture that such observers are aware of corpuscles in a way that fails to hold for wavefunctions. Of course, there is an obvious sense in which the corpuscles are also "hidden," and Dürr et al. emphasized in 1992 (Dürr et al. 1993) that the only time we can have sure knowledge of the configuration of corpuscles is "when we ourselves are part of the system." But how exactly is this supposed to work? Stone correctly pointed out in 1994 (Stone 1994) that this claim certainly fails if our knowledge is based on measurements which one part of our brain makes on another... (Brown and Wallace2005, p. 534) 849
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In support of premise 5 (the lack of a simple correlation between brain states and particle positions), Brown and Wallace point out: 859
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Suppose we accept that it is the [particle positions] that determine the outcome of the measurement. Is it trivial that the observer will confirm this result when he or she "looks at the apparatus"? No, though one reason for the nontriviality of the issue has only become clear relatively recently. The striking discovery 861
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865 in 1992 of the possibility (in principle) of “fooling” a detector in de Broglie–
866 Bohm theory (Englert et al.1992, Dewdney et al.1993, Hiley et al.2000, Brown
867 et al.1995) should warn us that it cannot be a mere definitional matter within
868 the theory that the perceived measurement result corresponds to the “outcome”
869 selected by the hidden corpuscles (Brown and Wallace2005, p. 523).

870 As premise 6 indicates, Bohmians might respond to this problem by appealing the
871 theory of decoherence. Decoherence involves considering how the action of two sys-
872 tems (thought of as the measuring apparatus and the object under study) on the wider
873 environment can enable them to become approximately classical in their relation to
874 each other, in such a way that they can be assigned stable properties (such as location)
875 that evolve in roughly the way prescribed by classical, pre-quantum physics.

876 However, it is not at all clear that decoherence will work in the intended way in a
877 Bohmian setting. Sanz and Borondo (Sanz and Borondo 2003) studied the double-slit
878 experiment in the framework of Bohmian mechanics and in the presence of deco-
879 herence and showed that even when coherence is fully lost, and thus interference is
880 absent, nonlocal quantum correlations remain that influence the dynamics of the par-
881 ticles in the Bohm theory, demonstrating that in this example decoherence does not
882 suffice to achieve the classical limit in Bohmian mechanics. See also (Schlosshauer
883 2005, 1298).

884 Is this problem of perceiving *pointer settings* any greater for the Bohmians than
885 it was in classical, Newton-Maxwell physics? Yes, it is, precisely because of the
886 radically non-local character of Bohmian dynamics. All distant bodies in Newto-
887 nian mechanics have a negligible influence on local phenomena, an influence that
888 decreases proportionally to the square of the distance. This is not the case in Bohmian
889 mechanics. There is, therefore, real grounds for doubting whether we can reliably
890 detect the actual positions of Bohmian particles, contrary to principle *E1*.

891 **7.5 The grw/objective collapse programme**

892 The hylomorphic interpretations of quantum mechanics have several advantages over
893 GRW and other non-hylomorphic objective collapse theories. First, hylomorphism
894 does not require speculation about some as-yet-unknown mechanism by which quan-
895 tum waves collapse into precise states. Consequently, hylomorphists can give a much
896 simpler account of the internal dynamics of the quantum world: the quantum world
897 proceeds without exception according to the dynamics of the Schrödinger equation.
898 Instead of postulating some unknown quantum trigger of wave collapse events, the
899 hylomorphic pluralist simply relies on our actual practice of using instruments with
900 classical features to precipitate precise measurement events. For hylomorphic plural-
901 ists, to learn more about how quantum waves collapse is simply to learn more about
902 macroscopic and mesoscopic systems themselves—to learn more chemistry and ther-
903 modynamics and biology. This is in fact the direction taken by generalized quantum
904 mechanics (as I described in Section 5).

905 In addition, the hylomorphist can take the objects of the ‘mesoscopic’ world
906 (including molecules and cellular structures) as persisting in stable states through
907 time, while the objective collapse view has to be combined with a further account of

the ontology of the macroscopic world. For example, if the GRW theory combined with John Bell's "flash ontology" (Bell1987, Maudlin2011, pp. 23–57), in which the macroscopic world consists of a number of widely separated and intermittent "flashes" (like the blinking of a swarm of fireflies), with each flash representing a wavepacket collapse. However, the Bell flash ontology can only provide a relatively small number of "flashes" of determinacy, too small a number to ground the existence of stable molecules and organisms:

The alternative version of GRW theory is the matter density interpretation. On this view, objective collapses result in relatively dense concentrations of expected mass in spacetime regions that resemble the objects of our classical world. The matter density interpretation shares with Bohmian theory the problem of verifying the reliability of our sense perception, and for similar reasons (both theories involve a high degree of causal non-locality). As Schlosshauer has pointed out, decoherence is of relatively little help to objective collapse theories (Schlosshauer 2005, pp. 1293–6).

In addition, as Alexander Pruss has recently argued (Pruss 2015), non-hylomorphic objective collapse theories face a problem with respect to the epistemological constraint E2, the reliability of memory. GRW is not really a single theory but a family of theories. The family has a single free parameter, which we can call (following Pruss) f , the *hitting frequency*. The hitting frequency gives us the probability of the collapse of any system of entangled particles, as a function of the total mass of those particles. We can put an upper bound on the hitting frequency—if f were too high, then we would never observe the kind of entanglement that is characteristic of the quantum realm. However, this experimental data puts no lower bound on the f . The frequency could be so low that it is very unlikely that any system should ever collapse. The argument against such a low frequency has to be philosophical and phenomenological rather than scientific: if the frequency were that low, human observations would never have definite or delimited outcomes, contrary to our experience.

Pruss suggests that we take such low frequencies seriously:

But imagine f is so low that typically a collapse in something the size of my immediate environment occurs only every hour. On its face this is ruled out by my memory of the past five minutes. But suppose, as seems reasonable on GRW, that consciousness occurs only when there is no superposition of brain states that constitute consciousness. Nonetheless, even when consciousness does not occur, my brain states will be evolving in superposition, and when they collapse they will give rise to conscious false memories of having had conscious states over the past period of time. We thus have no way of empirically ruling out such low values of f .

In other words, the proponents of GRW can rule out such low hitting frequencies by assuming (without argument) that our memories are veridical. However, the GRW family of theories, if true, would give us good reason to doubt that veridicality. If GRW were true and the hitting frequency were low, my *present experience* would be exactly the same. I could know that I have just now experienced a collapse of the wave function, but I could not have any confidence that any of my apparent memories of precise observations in the past are veridical. It isn't just that proponents of GRW are,

953 like all of us, subject to Cartesian doubts. It's rather that the GRW program provides
954 positive support to the skeptic's worries. If the hitting frequency is low enough, my
955 memories are radically unreliable as manifestations of the actual past. Some degree
956 of reliability is a condition of knowledge.

957 The defenders of GRW might object to this reduction to skepticism by arguing that
958 it is legitimate for them to take into account the need to secure the reliability of our
959 memory in fixing the value of the hitting frequency parameter. Why can't they simply
960 build a sufficiently high hitting frequency into their theory as a way of blocking the
961 argument for skepticism?

962 I have two responses. First, since f is a free parameter of the theory, the only
963 legitimate way to settle its value is empirically. However, its value cannot be settled
964 empirically without presuming (at least implicitly) that our memories are indeed reli-
965 able. Hence, it would be viciously circular to set the frequency high enough to ensure
966 the reliability of our memory. In contrast, the hylomorphist treats the reliability of
967 our memory as a fundamental fact about the human form, with no free parameters
968 whose value-determination requires empirical input.

969 Second, the GRW theorist is vulnerable to epistemic defeat, along the lines
970 developed by Alvin Plantinga (Plantinga1993, Plantinga2003, Plantinga2011). In the
971 absence of any physical or metaphysical constraints on the value of f , we have to
972 take seriously the possibility that the value of f might be extremely low. We know
973 that our memory is very unreliable, on the assumption that f is low (most of our
974 apparent memories are illusory). In that situation, we cannot appeal to our memory
975 of the past to verify the reliability of our memory without obvious vicious circular-
976 ity. Thus, we cannot justify continued rational belief in the reliability of our memory,
977 given the real possibility of an undercutting defeater which cannot itself be defeated.

978 In contrast, there is no similar consideration forcing the hylomorphist to recognize
979 any possibility of the unreliability of our powers of memory.

980 Finally, even if we were to grant that the hitting frequency is so low that such false
981 memories would be extremely unlikely, this is not sufficient for our memory-based
982 beliefs to constitute knowledge. A very high probability of truth is not sufficient
983 for knowledge, as the famous *lottery paradox* illustrated. I can know that the prob-
984 ability of each ticket's winning is extremely low—in a hypothetical lottery with an
985 astronomical number of tickets, fantastically low. However, such a low probability
986 of falsity is not sufficient to give us knowledge of truth, since if I could know that
987 each ticket is a loser, I could also know that they all are, which in fact I know to be
988 false. What's needed for knowledge is the exercise of some cognitive power which,
989 if exercised in normal circumstances and without external interference, guarantees
990 absolutely the truth of the belief formed. Given GRW without hylomorphic powers,
991 our memory-based beliefs can never meet that standard.

992 Therefore, GRW theories and other objective collapse theories fail epistemological
993 constraint E2.

994 GRW theories also fail constraint E1, perception, for reasons noted by David
995 Albert and Lev Vaidman (Albert and Vaidman 1989) and (?Albert1990). The
996 human visual system is quite sensitive to small numbers of photons—as few as six or

seven suffice. However, such a small collection of photons has a vanishingly small probability of inducing a wavefunction collapse under GRW models. Aicardi et al. (Aicardi et al.) responded by arguing that the movements of ions in the human nervous systems that correspond to the apparent perception of photons is sufficient to guarantee a collapse with high probability within the time frame of conscious perception. However, this is not sufficient to satisfy E1, since it means that almost all of our visual perceptions are factually inaccurate. They represent events occurring in our environment, events that are ontologically independent of the movement of ions in our optic nerves and brains. If GRW is correct, however, what we see when we see something is actually an event occurring within our own nervous systems. There was no corresponding external event consisting of the emission of a localized photon that we were able to detect. Once again, GRW can *save the phenomena* but only at the expense of undermining human knowledge.

8 Conclusion

Power ontology provides us with a metaphysical framework that is sufficiently flexible to accommodate fundamental modes of causation at the level of thermodynamics, chemistry, and solid-state physics. By doing so, we can circumvent the usual measurement problem, which presupposes that an exhaustive description of the world at a fundamental level can be given in terms of pioneer quantum mechanics, with no non-trivial center of classical properties.

Additional work needs to be done in exploring the relationship between a purely quantal description of particles (taken either individually or as definite pluralities of discrete entities) and the metaphysically more fundamental level of substances and their causal powers. In particular, should we assume that there is a quantum wavefunction that embraces all the particles of the world, simultaneously characterizing the quantum potentialities of all substances, or should we suppose instead that quantum wavefunctions are always local and contingent affairs, part of what Nancy Cartwright has described as a *dappled world*? (Cartwright 1999) The hylomorphic view can be developed in either direction. If we assume a global wavefunction, then we get the *traveling forms* interpretation of Alexander Pruss, in which substantial forms of interacting substances induce global collapses of the wavefunction. (Pruss 2018) The dappled world alternative has been developed by William Simpson in his dissertation (Simpson 2020), and it is that model that is tacitly presupposed by Primas’s model of collapse. It also underlies recent work by Barbara Drossel and George Ellis.(Drossel and Ellis 2018)

This issue corresponds to a further question about the extent of entanglement in nature. The global wavefunction picture would suggest that entanglement is pervasive in nature, arising with the Big Bang and never fully disappearing. On the dappled world picture, entanglement occurs only under special circumstances, when complex systems are prepared in a way that is isolated from the surrounding environment. Local collapses destroy these fragile entanglements.

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